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Teomical Report Re. 14

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August 1155

Department of (hemistry Tiliseis Institute of Technology 5500 S. Federal Street Chicago 16, Illinois U. S. Navy
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### CHEMICAL THERMODYNAMICS OF MATERIALS AT HIGH TEMPERATURES

## Technical Report No. 14

# SOLID SOLUTION EQUILIBRIA IN THE TERMARY SYSTEM ZIRCONIOM - OZYGEN - HYDROGEN

by:

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August 1953

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Chicago 16, Illinois

## SOLID SOLUTION EQUILIBRIA IN THE TERNARY SYSTEM ZIRCONIUII - OXYGEN - HYDROGEN

By R. K. Edwards, P. Levesque, 1,2 and D. Cubicciotti3

- (1) Based on part of a thesis submitted by Pascal Levesque to the Graduate School of Illinois Institute of Technology in partial fulfillment of the requirements for the degree of Doctor of Philosophy, June, 1953.
- (2) Presently at Metallurgy Department, Sylvania Electric Corporation, Ipswich, Massachusetts.
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#### ABSTRACT

The solubility of hydrogen in the solid solution of the ternary system zirconium-oxygen-hydrogen has been determined as a function of temperature, composition, and equilibrium hydrogen pressure. The data have been correlated with those of the related binary systems and some previous work in the ternary system, and the solid state chemistry has been represented by an extensive provisional partial ternary phase diagram.

A new single-phase region,  $\mathcal{E}$ , lying between the  $\alpha$ -zirconium and  $\mathcal{E}$  (approximately  $ZrH_2$ ) phases has been established as belonging both in the binary and ternary systems.

The hydrogen saturation boundary of the  $\alpha$ -phase shows that initially hydrogen may be replaced by oxygen on a three-to-one basis in the  $\alpha$ -phase. The hydrogen saturation boundary of the S-phase indicates that hydrogen may be replaced by oxygen on a three-and-one-half to one basis in the S-phase.

\* \* \* \* \* \*

#### INTRODUCTION

The solid solubility of oxygen in zirconium metal has long been considered to be very high.4,5 A recent thorough phase diagram study in

- (4) J. H. de Boer and J. D. Fast, Z. anorg. Chem., <u>148</u>, 345 (1925).
- (5) D. Cubicciotti, J. Am. Chem. Soc., 73, 2032 (1951).

the zirconium-oxygen system by Domagala and McPherson<sup>6</sup> has set the solubility

(6) R. F. Domagala and D. J. McPherson, "Phase Diagrams of Zirconium-Base Binary Alloys, Report No. 10, The Zirconium-Oxygen System, Report 4 - Summary COO-181," Armour Research Foundation, Chicago, Illinois, March 31, 1953.

at 29 atom per cent oxygen in the temperature range 600 to 900°. We have previously presented a partial phase diagram for the zirconium-hydrogen system and have indicated a primary solid solubility for hydrogen in zirconium of from 49 atom per cent at 700° to 52 atom per cent at 875°.7 Hall, Martin,

(7) R. K. Edwards, P. Levesque, and D. Cubicciotti, "Solid Solution Equilibria in the Zirconium-Hydrogen System," Tech. Report No. 13, U. S. Navy, Office of Naval Research, Contract N7-ONR-329, Task Order II, August (1953).

and Rees, who studied the effect of oxygen on the solubility of hydrogen in

(8) (Mrs.) M. N. A. Hall, S. L. H. Martin, and A. L. G. Rees, Trans. Faraday Soc., 41, 306 (1945).

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(9) D. P. Smith, "Hydrogen in Metals," The University of Chicago Press, Chicago, 1948.

behavior with regard to his "rift theory" as the mode of occlusion of hydrogen in metals.

It occurred to us that the observed "equivalence" might be nothing more than a manifestation of the conversion of  $ZrH_2$  to  $ZrO_2$  on addition of oxygen since these are the "saturation"phases with respect to the two components, hydrogen and oxygen. On the other hand the wide solid solubility, discussed above, of both hydrogen and oxygen in zirconium would lead to a prediction of extensive ternary solubility. One might envisage this ternary solution as a metallic zirconium matrix lattice with the interstices variously filled

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with the small hydrogen and oxygen particles. The idealized hexagonal close-packed lattice can provide both octahedral and tetrahedral interstices. We might expect some rather striking manifestations as oxygen replaces hydrogen in view of their conspicuous differences in size and in electron affinities. We have previously mentioned the work of Coehn and coworkers which showed hydrogen in palladium solid solution to migrate with passage of an electrical current as though it were a positively charged particle. In contrast, de Boer and Fast and Almutt and Scheer have

(10) D. B. Alnutt and C. L. Scheer, Trans. Electrochem. Soc., <u>88</u>, 195, (1945).

shown that oxygen in zirconium migrates as a negatively charged particle.

The present investigation was conducted by the method of measuring the equilibrium hydrogen pressure over ternary solid solutions as a function of composition and temperature. The composition was varied by successive additions of measured volumes of hydrogen to an initially homogeneous zirconium—oxygen solid solution in which the composition ratio of zirconium to oxygen remained constant.

#### EXPERIMENTAL

Each zirconium—oxygen solid solution specimen was prepared by controlled oxidation of about 100 mg. of 0.015—inch diameter zirconium wire, followed by equilibration of the resulting metal and oxide mixture for twenty minutes in vacuum at a temperature approximately 200° below the melting point of the particular composition. Equilibration heating was by electrical resistance of the wire in vacuum. The apparatus is shown in Figure 1, and the procedure has been previously described. The homogeneous nature of each preparation was carefully established by x—ray diffraction examination before it was considered a suitable specimen for study. After the composition was established by weight difference on oxygen uptake as observed on an Ainsworth semi-micro balance, a middle section of the wire was cut out for use in the hydrogen solubility determination.

Numerous preliminary attempts at equilibration by holding a preparation at 1200° for twelve hours did not always result in a homogeneous single phase solution, and the procedure was abandoned as unreliable. Hall, Martin, and Rees reported the latter procedure to be satisfactory at 950° in their work.

#### MATERIALS

Zirconium wire of 0.015-inch diameter obtained from the Foote Mineral Company was used. The analysis has been given previously. Commercial 99.8 per cent pure tank oxygen was used without additional purification.

#### PROCEDURE

The apparatus and procedure for determining the equilibrium hydrogen pressure as a function of temperature and composition of the ternary material were those we had used in the binary zirconium-hydrogen study, and the same limits of error are assumed. Families of isotherms were determined for samples of initial oxygen compositions of 6.3, 10.8, 18.6, 24.7, and 29.6 atom per cent in zirconium, and these values correspond to atom ratios, R, of oxygen to zirconium of 0.068, 0.121, 0.229, 0.328, and 0.420, respectively.

#### RESULTS

The experimental pressure-composition isotherms are shown on two different coordinate bases. Figures 2, 3, 4, 5, and 6 show families of isotherms of equilibrium hydrogen pressure as a function of the volume of hydrogen gas taken up per gram of zirconium metal in a sample of given R value. Figures 7, 8, 9, 10, and 11 show corresponding plots of the logarithm of the hydrogen pressure versus the logarithm of the atom per cent hydrogen composition. Advantage was taken of the reasonably linear nature of the extrapolations in the logarithmic plots to obtain the phase field boundary compositions plotted in Figures 12, 13, 14, 15, and 16, and listed in Table 1. In Figures 2 and 3, one sees the unmistakable appearance of a second composition region of pressure arrest in addition to the one corresponding to the arrest in the binary system. The data shown on the 800° isotherm in Figure 2 are the results of four different determinations and amply establish the existence of the second arrest.

On the other hand, Figures 4, 5, and 6 show only wide composition ranges of a single pressure arrest. Values of equilibrium hydrogen pressure as a function of temperature at several significant fixed compositions have been taken from the smoothed curves of Figures 2, 3, 4, 5, and 6 and the data are tabulated in Table 2 and plotted in Figures 17, 18, 19, 20, and 21 with the logarithm of the equilibrium hydrogen pressure as a function of the reciprocal temperature. From the plots partial molal heats of solution

of molecular hydrogen in the ternary solid solutions were obtained, and these values are also given in Table 2.

A summarizing correlation is given in Figure 22 where an isothermal section of a provisional ternary phase diagram is shown in part. The data for the binary zirconium-hydrogen system boundaries were taken from our previous paper. The data for the binary zirconium-exygen system boundaries were taken from Domagala and McPherson. The data of the composition region of R value greater than 0.420 we have obtained from replotting the data of the isotherms of Hall, Martin, and Rees, in the manner of the handling of our own data.

#### DISCUSSION

The solid state chemistry of the zirconium-oxygen-hydrogen system in the temperature range 600 to  $900^{\circ}$  and at hydrogen pressures of less than 760 mm. of Hg can be shown by a quasi-ternary phase diagram of  $Zr-ZrO_2-ZrH_2$  as is presented in Figure 22. It is to be borne in mind that actually somewhat less than the hydrogen composition corresponding to  $ZrH_2$  is reached under the given temperature and pressure conditions. Nomenclature of phase fields is that previously adopted.

A three-component system in the presence always of a gas phase will yield invariant isothermal pressure behavior with changing composition if three solid phases are present. In the following discussion the presence of a gas phase is to be understood and only the solid phases are explicitly discussed. The two-phase regions, however, may appear to show equivalent behavior if the change in composition happens to be effected along a tie-line connecting phase compositions of equal fugacity. As compositions are continuously and isothermally changed along a path traversing first a single phase, second a two-phase, and third a three-phase region, one would expect the following general behavior. The pressure isotherm would change continuously in the one-phase region and also in the two-phase region, with a change in slope at the join of these regions, and it would finally change in slope at the boundary and remain invariant throughout the three-phase region.

An interesting consequence of the natural phase relations and the manner in which composition changes were effected in the present study is that for each material of a given R value, pressure-composition isotherms clearly indicate that the composition traverses were essentially along

tie-lines in the two-phase regions. That this is not three-phase behavior can be seen by noting that pressures do actually vary in the regions indicated in Figure 22 to be two-phase regions. A three-dimensional plot having the isothermal ternary diagram as a base and hydrogen fugacity as the vertical ordinate would readily demonstrate this point. We have indicated the observed behavior in the ternary diagram by showing the two-phase regions shaded with lines presumed to closely approximate lines of constant fugacity. Induction of the nature of the phase diagram was based on the above, and on the additional requirements that adjoining single-phase regions must be separated by a two-phase field and adjoining two-phase fields must be separated by a three-phase field. Room temperature x-ray diffraction patterns were taken on all of the fully hydrogenated specimens and in all cases the E-phase? accounted for all diffraction lines. In particular, and somewhat to our surprise at first, none of the ZrO<sub>2</sub> diffraction pattern was found for any of these specimens. The diagram given is consistent with this observation.

A preparation of R = 0.121, of such a hydrogen composition as to place it just inside the 6-region on the sirconium-rich side, was examined by x-ray diffraction at room temperature. The relative intensities and positions of the lines of the pattern, except for one line, could be reasonably well correlated with patterns of a-zirconium, assuming appreciable distortion of the lattice. However, the one additional line of medium to strong intensity and low Brigg angle was unmistakably present. The broadness of the lines, like that generally found in the patterns taken on materials in the binary sirconium-hydrogen system, precluded more detailed structural analysis.

The partial molal heats of solution of molecular hydrogen in the ternary solid solutions are found to be the same as in the binary sirconium-hydrogen system? for the  $\alpha$  phase and  $(\alpha + \delta)$  two-phase fields for the materials of R values 0.068 and 0.121. For the latter, the  $(\delta + \epsilon)$  two-phase field yielded  $\Delta H = -29.7$  Kcal per mole as did the  $\alpha$  phase and  $(\alpha + \epsilon)$  two-phase fields for all higher R values. It is to be added that the significance of the  $\Delta H$  values for the two-phase regions may be questioned since their composition boundaries are not constant, as is discussed briefly below.

The temperature dependence of the various phase field boundaries may be seen in Figures 12, 15, 14, 15, and 16. One may note a tendency toward closure of the  $(a + \delta)$  and  $(\delta + E)$  fields at higher temperatures.

We have not discussed the presence of the  $\beta$ -phase region in this ternary study since from the known binary zirconium-oxygen system behavior<sup>6</sup> it is most probable that the region would have been displaced to temperatures higher than those used in this study.

One may refer to the x-ray diffraction patterns by Sidhu and McGuire11

(11) S. S. Sidhu and J. C. McGuire, J. App. Phys., 25, 1257 (1952).

for an idea of the magnitude and nature of the changes in the diffraction pattern as hydrogen is added to hafnium metal. 12

(12) The quality of the patterns obtained with hafnium is considerably superior to that obtained with zirconium.

We urge a reasonable caution in the use of the diagram we present and wish to note in particular the uncertainty in the actual apex positions of the three-phase rields. The extent of the uncertainty may be easily inferred on noting the amount and nature of the extrapolation of boundaries to these apex positions. The extrapolation to the  $(\delta + \Xi)$ two-phase field boundaries leads to boundary compositions of about 61 and 62 atom per cent hydrogen for the binary zirconium-hydrogen system. We suppose that the value 62 atom per cent given for the hydrogen-rich boundary is probably quite accurate since the linear extrapolation nicely parallels the isothermal composition isobers, one of which is shown for the pressure 760 mm. of Hg by the dotted line in the E-region of Figure 22. On the other hand, a strictly linear extrapolation of the data for the sirconiumrich ( $\delta + E$ ) border would have led to a composition only negligibly lower than 62 atom per cent so that it may well imply an extremely narrow corresponding region in the binary zirconium-hydrogen system. Furthermore, if the boundary were taken as linear through the data points plotted, it would imply the relationship that the hydrogen saturation limit of the & -phase decreases in a manner which shows oxygen to replace hydrogen on a three-and-one-balf to one atom basis. The rather extensive linear portion of the hydrogen-rich c-phase saturation boundary, similarly considered, implies that oxygen will replace hydrogen on a three-to-one atom basis.

In the a-zirconium hexagonal close-packed structure, one would suppose that the oxygen atoms would fill the octahedral interstices—the largest holes available. Apparently they cannot do so completely or the solubility limit of oxygen in zirconium in the binary zirconium-oxygen system would correspond to 50 atom per cent. The actual limit of 29 atom per cent shows that about two out of every five available octahedral interstices in a-zirconium become filled at saturation with oxygen.

On the other hand, hydrogen may be either in tetrahedral interstices or in the octahedral interstices. If it is in the latter, the hydrogen binary saturation composition of 50 atom per cent in the zirconium-hydrogen system would indicate that all available octahedral interstices are filled. It appears then that each oxygen atom which enters the lattice denies more of the interstitial sites to hydrogen than it does to oxygen itself. This would seem to indicate a form of repulsion through general lattice strain and to rule out any special stabilization one might have anticipated to arise from proximate metallic hydrogen. Effects due to the difference in electron affinities between oxygen and hydrogen seem to be absent.

The dotted line in Figure 22 represents the contour of an isothermal composition isobar for a hydrogen pressure of 760 mm. of Hg and is therefore our value of the "saturation" composition as the term was used by Hall, Martin, and Rees.<sup>8</sup> It is seen that oxygen does replace hydrogen on approximately a one-to-one basis as supposed, but that this point is of trivial interest since "saturation" in that sense represents no phase field limit.

Table 1 Boundary Compositions of Solid Two-Phase Fields in the

Zirconium-Oxygen-Hydrogen System

(Compositions in Atom % Hydrogen)

R		<del></del>	T	EMPERA	rures	-	
			600°	700°	750°	8 <b>00°</b>	850°
0.068	(two two-phase regions)	First			42.8	42.9	45.7
				50.1	51.2	51.2	53.7
		Second		56.6	56.2	55.6	antonibe
			-	58.0	58.0	58.0	
0.121	(two two-phase regions)	First		34.8	<b>3</b> 5.9	36.7	39.1
				42.7	43.4	44.2	47.0
		Second	51.9	51.9	51.8	51.5	670KIRC#80
			55.6	55 <b>.3</b>	55,0	54.3	
0.229	(single two-phase field)		450,,,,,,,,,,	25.7	25.3	25.0	emperiodicality (in the contraction of the contract
			<b>49.</b> 0	40,4	49.6	49.6	e()=======
0.328	(single two-phase field)			15.0	17.1	17.1	ORNIA SIMPONO
			45.2	45.7	46.8	47.3	OCCUPATION .
0.420	(single two-phase field)		ann a state of	9.6	12.0	12.0	macediffectatio
			<b>3</b> 8.9	38.9	38.9	39.3	RacMEDIED

Table 2 Equilibrium Hydrogen Pressure as a Function of Composition in the

Ternary Zirconium-Oxygen-Hydrogen System

(Pressures in mm. of Hg.)

ρ	Phase	Hydrogen Composition	AH of			TEMPERATURES	rures		
•	Region	Atom %	Kcal./mole	<b>6</b> 009	200	750	800	850	<b>9</b> 006
890	В	31.4	-39.6	Carponer		10	27	9	125
	ಕ	37.9	-39.6		6	24	65	145	270
	+	47.8	49.6	coleman	13	22	165	480	
	S + E	57.2	GENERAL STREET	15	140	230	472		
	W	60,4		8	450	099	Canada	-	
0.121	8	30.4	-39.6		ogentuise	18	48	300	212
	<b>3</b> + 5	42.1	49.6	(i)	13	45	25	455	
	4	48.6	CARCONIC		23	120	342	605	accompany
	S+ E	<b>22.</b> 8	2.6.7	43	165	310	595	C	Calendaria
	W	28.0	Constitution	280	260	775			oct scars
0.229	೪	11.7	-29.7		80	40	80	145	250
	ช	16.6	-29.7		46	95	<u>00</u>	325	210
	₩ + !	29.8	-29,7	32	135	200	515	encocine encocine	COCOCOCO
	لد	52.0	COLUMN TO SERVICE	285	485	648		Constant of	GLOCI
0.528	ಆ	5,8	-29.7	6.0E.X.2E.3	12	24	49	85	155
	ម	10.9	-29.7	10	<b>ب</b>	20	160	230 230	440
	W + 8	35, 5	-29.7	30	ᅋ	275	465	00.11.12	C42.40
	W	49.5		205	440	650			
0.420	ರ	7.9	-29.7	13	ន	100	220	375	contro
	₩ + 8	25,6	-29.7	22	92	303	390	COLUMN	
	W	42.7	(35.4.4.25)	110	285	450	099	cocco	
							James Contraction	Caronovious and	

Table 3 Equilibrium Hydrogen Pressure as a Function of Composition and Temperature

(Experimental data for the oxygen-zirconium atom ratio, R = 0.068)

Pres-	Atom \$	Pres-	Atom \$	Pres-	Atom %
sure	Hydro-	sure	Hydro-	sure	Hydro-
mm.Hg	gen	mm.Hg	gen	mm.Hg	gen
Run 165	-600°	Run 166	<u>-750</u> °	Run 157	-850°
0	38.7	177	54.3	22	20.0
12	55.0	257	56.1	53	30.4
60	60.2	<u>Run 158</u>	3 <b>–</b> 800°	102	36.2
136	60.6		1	167	38.9
255	61.0	49	<b>35.</b> 8 <b>41.3</b>	198 232	39.5 40.6
407	61.9	123 164	48.1	312	42.6
592 726	62.5 62.8	201	51.5	376	43.8
1	•	· ·		443	45.0
Run 160	<u>-700</u> °	<u>Run 161</u>	<u>-800</u>	473	47.2
14	43.5	14	23.4	502	54.2
20	51.1	93	40.0	586	54.9
89	55.5	147	42.6	720	56.4
142	57.0	289	53.6	Run 159	
288	59.4	714	59.3	679	55.5
Run 170	<u>-700</u> °	Run 163	5 <u>800</u> °	Run 169	9 <b>−</b> 850°
14	43.3	241	52.4	108	35.5
49	53.8	<b>37</b> 8	54.5	205	39.6
145	58.1	618	58.6	310	42.3
167	58.3	724	59.3	Run 159	9-900°
277	59.4	<u>Run</u> 16	, 7 <b>⇒</b> 800°	•	
383	60.2	[ ]		26	17.2
51.7	60.6	38	33.6	63	25.2
624	61.6	119	41.1	123 184	31.0 34.9
731	62.0	190 360	51.9 54.8	253	37.3
Run 164	4-750°	405	55.6	317	40.0
9	27.3	460	56.2	396	42.0
27	38.0	474	58.2	521	44.6
48	41.6	51 ·	1	627	45.6
76	52.1	<u>Run 168</u>	<del>3-800</del>	734	
126	53.7	135	41.2	<u>Run 16</u>	
196	55.0	303	53.9		
316	58.0	41.5	55.3	324	39.2
474	59.5	439	55.8	417	42.3
476	59.1	460	56.1	523	43.7
5 <b>7</b> 8	60.1	478	57.5	629	45.2
712	60.8	500	58.0	1	
1		560	58.6	1	
ı		649	59.2		

table 4 Equilibrium Hydrogen Pressure as a Function of Composition and Temperature

(Experimental data for the oxygen-zirconium atom ratio, R = 0.121)

Pres- sure mm. Hg	Atom \$ Hydro— gen	Pres- sure mm.Hg	Atom \$ Hydro- gen	Pres- sure mm.Hg	Atom % Hydro— gen
Por 201	<u>=600</u> °	Run 203	<u>-750</u> °	Run 198	<u>-850</u> °
1 9 35 45 87 275 477 615 740 8m 200 5 13 21 44 95 146 175 220 332 521 598 743	39.0 50.4 53.0 56.6 57.1 58.0 58.9 59.3 59.6	12 30 45 64 122 197 287 311 348 603 743 Run 197 10 38 66 96 133 154 179 195 245 314 360 409 431 492 542 598 662 716 751	25.1 32.6 40.8 45.4 47.9 49.3 51.3 53.5 55.5 56.7 57.0	21 50 151 257 Run 199 16 105 336 391 440 462 498 541 614 725 Run 202 24 78 132 230 332 457 636 741	16.8 23.4 32.1 34.4 -850° 15.4 30.2 36.9 38.2 39.1 45.5 46.2 47.9 49.0 49.8

Table 5 Equilibrium Hydrogen Pressure as a Function of Composition and Temperature

(Experimental data for the oxygen-zirconium atom ratio, R = 0.229)

Pres- sure mm. Hg	Atom % Hydro- gen	Pres- sure mm.Hg	Atom % Hydro- gen	Pres- sure mm.Hg	Atom % Hydro- gen	Pres- sure mm.Hg	Atom % Hydro- gen
hun 17	-600°	<u>Run 175</u>	-750°	<u>Run 172</u>	2 <u>-850</u> °	Run 173	5800°
17 32 46 80 146 263 425 566 586 741  Run 179 599  Run 180 29 57 81 117 131 164 220 327 458 600 727  Run 181- 48 67 85 118 135	53.6 -700° 14.7 19.8 23.8 28.1 42.6 48.2 50.9 51.8 52.9 53.4 54.1	132 275 305 321 353 411 488 570 690 Run 176 50 108 166 248 Run 177 20 41 67 106 152 202 258	12.4 15.4 18.0 22.2	25 60 112 201 274 341 410 464 516 565 623 749 Run 171 48 114 194 301 410 543 656 752	7.2 9.9 11.6 13.6 14.5 16.4 18.0 19.8 20.2 20.1 21.2 23.2 -900° 8.3 10.1 12.3 15.6 17.5 17.9 18.7	38 97 157 210 280 333 412 462 515 Run 174 50 107 155 214 270 305 372 437 506 Run 178 438 505 505 535 576 638 731	9.8 12.7 14.4 15.8 19.0 19.2 22.8 24.8 26.4 -800 10.2 13.1 14.0 16.0 18.3 19.0 21.2 24.7 25.2

Table 6 Equilibrium Hydrogen Pressure as a Function of

Composition and Temperature

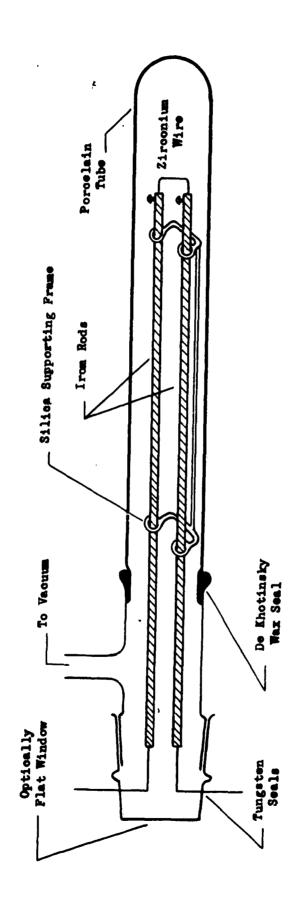
(Experimental data for the oxygen-zirconium atom ratio, R = 0.329)

Pres Atom % Pres Atom % Sure Hydromm.Hg gen mm.Hg gen	Pressure	Atom % Hydro- gen
Run 186–690° Run 183–750°	Run 188-	-850°
12   24.0   24   7.1   21   33.6   61   10.0   26   42.9   154   13.7   204   15.6   255   16.4   255   16.4   277   28.7   289   38.6   277   47.9   203   49.4   49.1   310   50.2   197   15.2   284   46.0   617   51.8   405   48.4   725   52.4   8un 185a-750°   8un 189-700°   570   49.1   37   10.2   650   49.6   61   13.5   737   50.0   89   18.2   107   35.1   128   43.0   128   43.0   128   44.2   251   48.1   373   48.7   373   49.1   526   49.9   685   50.5   29   5.4   62   8.5   137   10.2   199   11.0   300   12.9   400   16.1   441   18.2   454   32.4   474   41.9   723   48.5	73 175 255 352 464 554 619 727 Run 189- 37 71 176 278 378 483 569 655 752	5.6 7.1 9.2 11.6 14.8 16.4 18.4 20.1

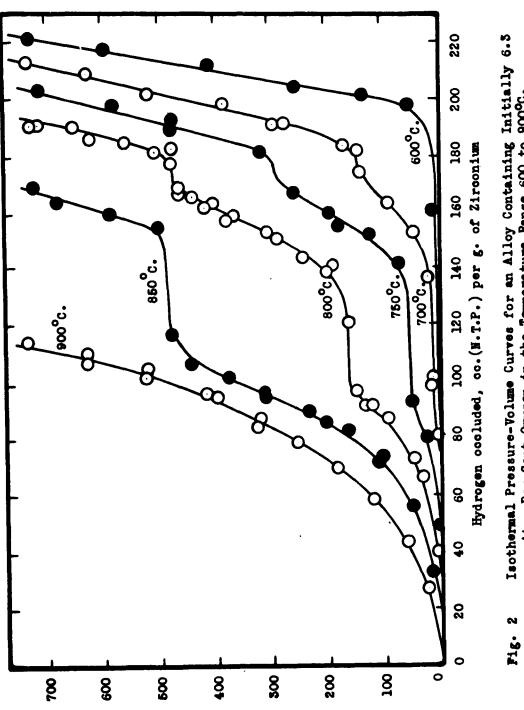
Table 7 Equilibrium Hydrogen Pressure as a Function of Composition and Temperature

(Experimental data for the oxygen-zirconium atom ratio, R = 0.420)

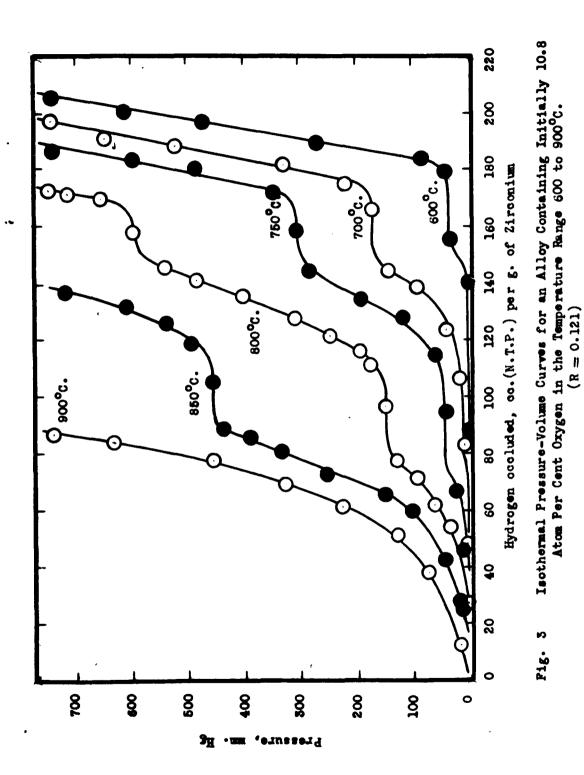
Pressure	Atom % Hydro- gen	Pres- sure mm.Hg	Atom % Hydro— gen	Pres- sure mm.Hg	Atom % Hydro- gen
Run 193 18 30 53 117 223 396 577 751 Run 192 32 75 211 337 493 599 750 Run 196	15.2 36.8 41.4 43.0 44.1 46.0 46.8 47.5 3-700° 6.8 10.7 41.2 43.5 44.8 45.5 46.6 3-700°	Run 193 57 130 Run 193 21 62 153 193 205 258 390 497 604 733 Run 193 29 89 190 308	6.4 8.7 5-750° 6.0 6.8 9.3 17.0 34.0 40.0 41.8 43.0 43.7 44.8	Run 19 35 148 268 379 389 442 532 641 739 Run 19 26 99 196 324 478 585 739	4.9 5.9 8.3 12.5 28.1 38.9 41.5 42.7 43.7
94 1 <b>3</b> 5	31.7 39.6				

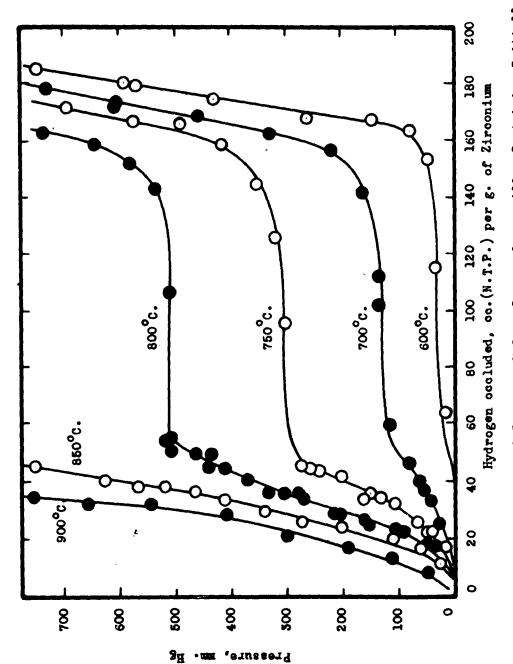


(Iron rods served as electrodes for conducting current through Zr wire) Apparatus for Preparing Zirconium-Oxygen Alloys. Fig. 1

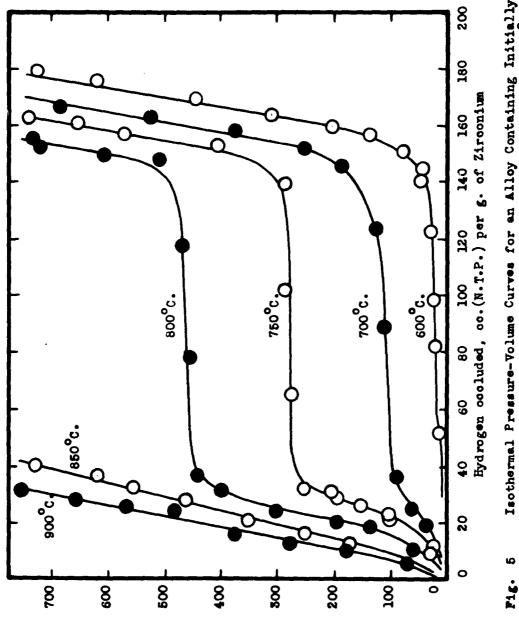


Isothermal Pressure-Volume Curves for an Alloy Containing Initially 6.5 Atom Per Cent Oxygen in the Temperature Range 600 to 900°C. (R = 0.068)

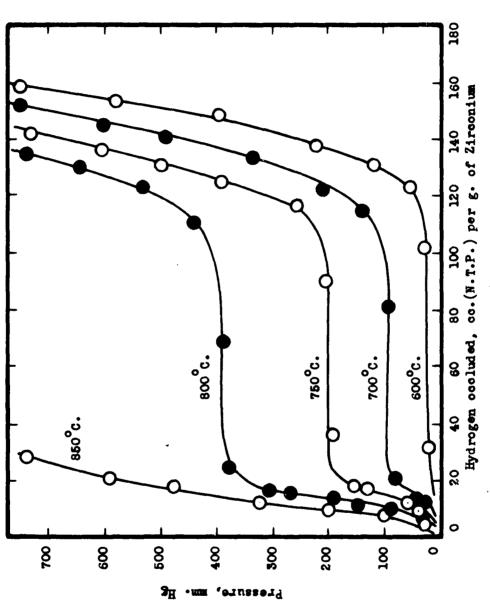




Isothermal Pressure-Volume Curves for an Alloy Containing Initially 18.6 Atom Per Cent Oxygen in the Temperature Range 600 to 900°C. (R = 0.229) F16. 4



Isothermal Pressure-Volume Curves for an Alloy Containing Initially 24.7 Atom Per Cent Oxygen in the Temperature Range 600 to  $900^{\circ}C$ . (R = 0.528)



Isothermal Pressure-Volume Curves for an Alloy Containing Initially 29.6 Atom Per Cent Oxygen in the Temperature Bange 600 to 850°C. (R = 0.420) Fig. 6

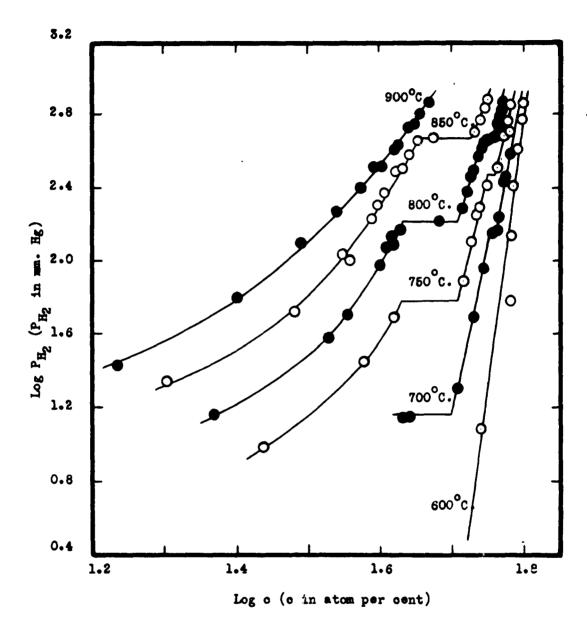


Fig. 7 Isothermal Pressure-Composition Logarithmic Plots for an Alloy Containing Initially 6.5 Atom Per Cent Oxygen.
(R = 0.068)

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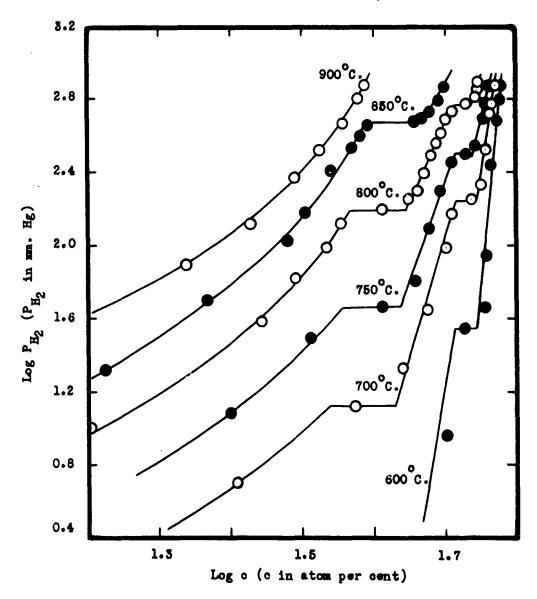
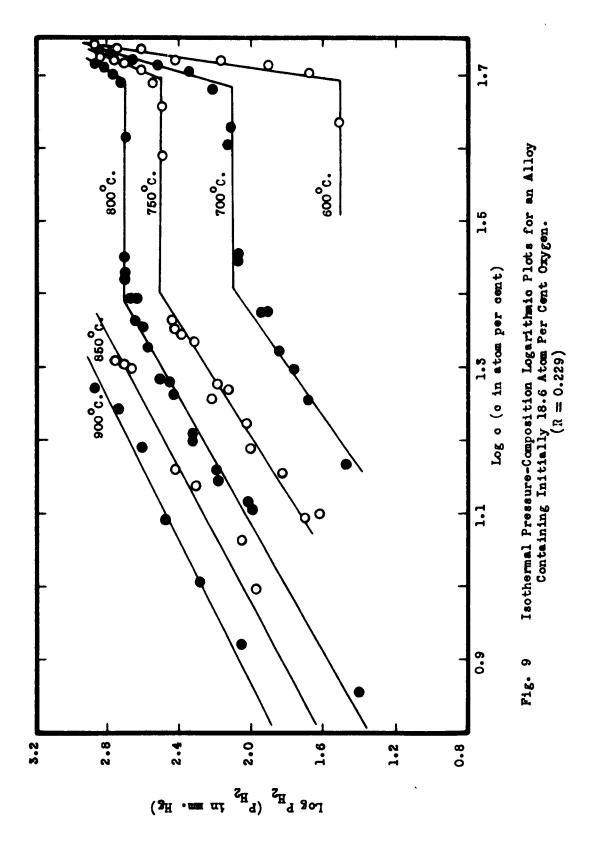
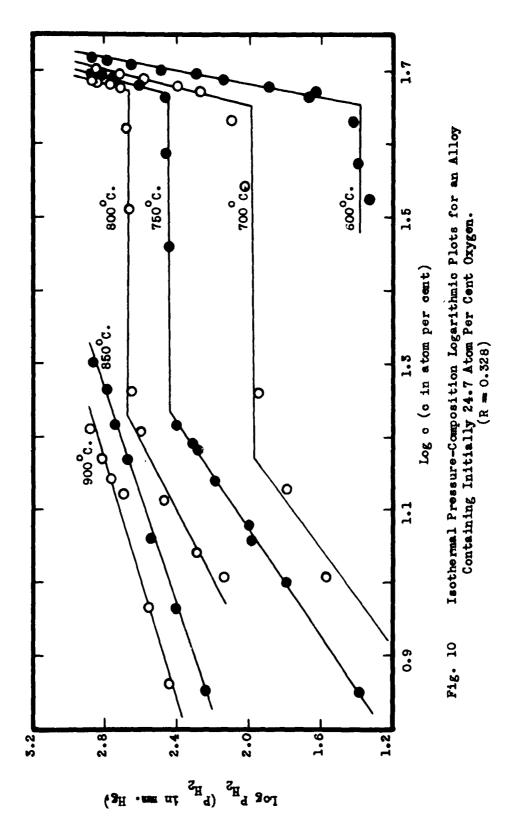
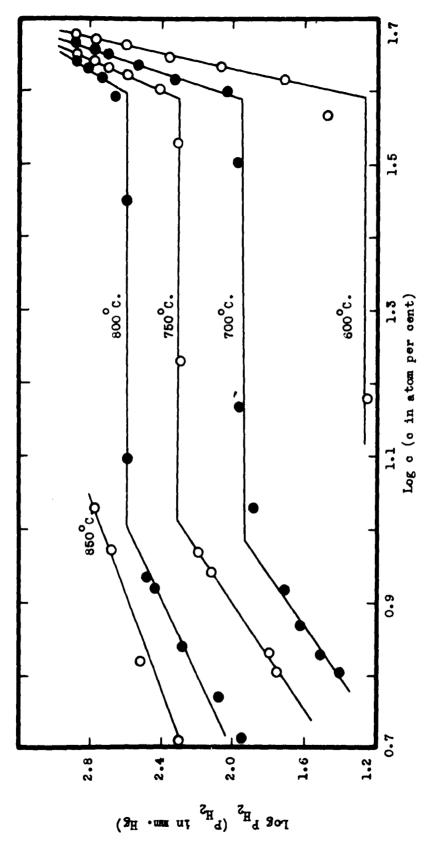


Fig. 8 Isothermal Pressure-Composition Logarithmic Plots for an Alloy Containing Initially 10.8 Atom Per Cent Oxygen.
(R = 0.121)



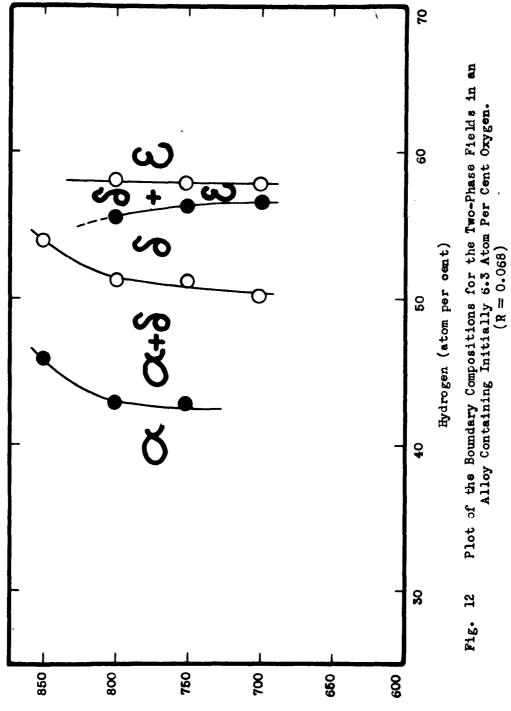




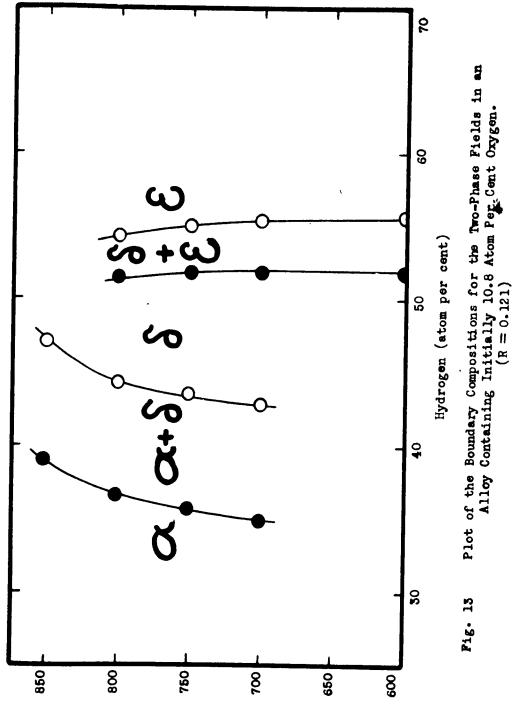
Isothermal Pressure-Composition Logarithmic Plots for an Alloy Containing Initially 29.6 Atom Per Cent Oxygen. Fig. 11

(R = 0.420)

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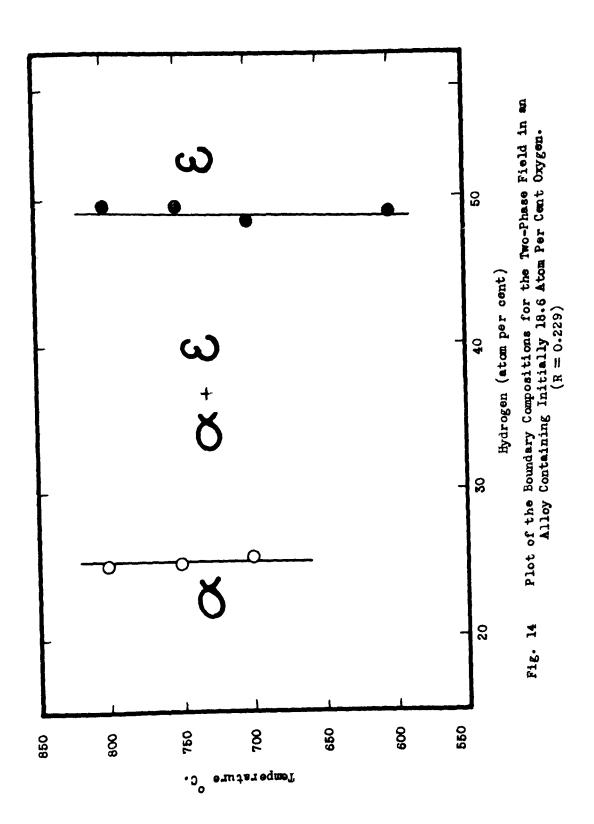


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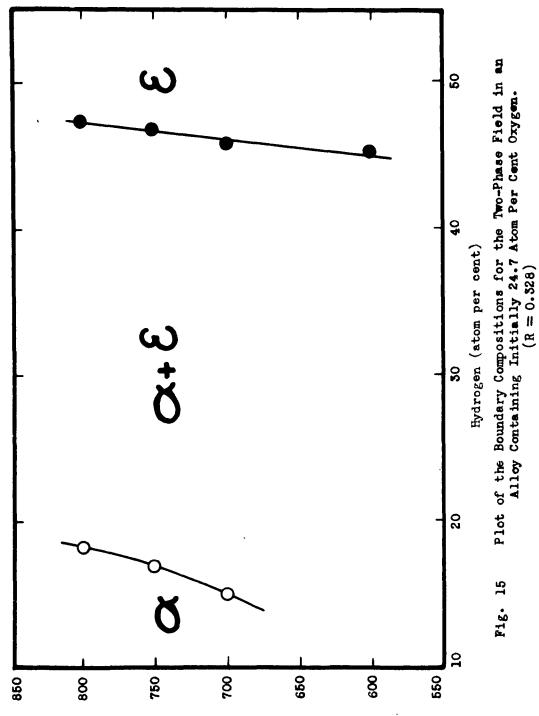
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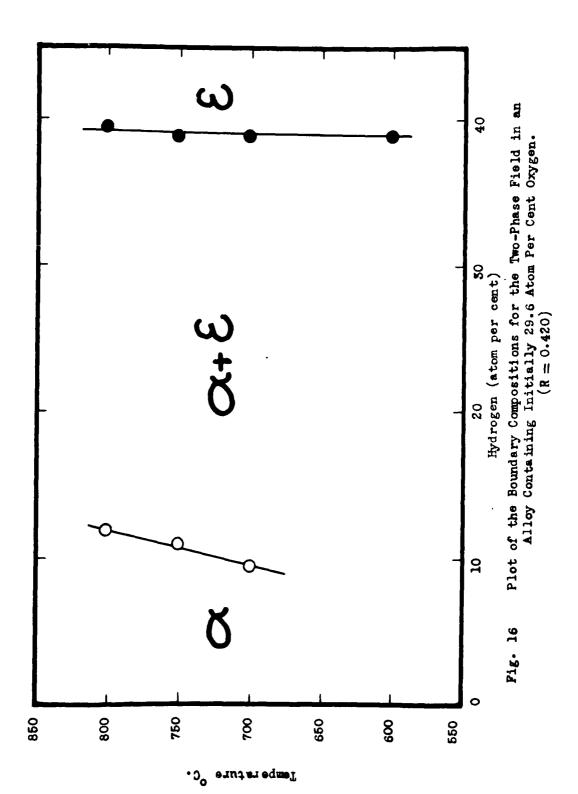
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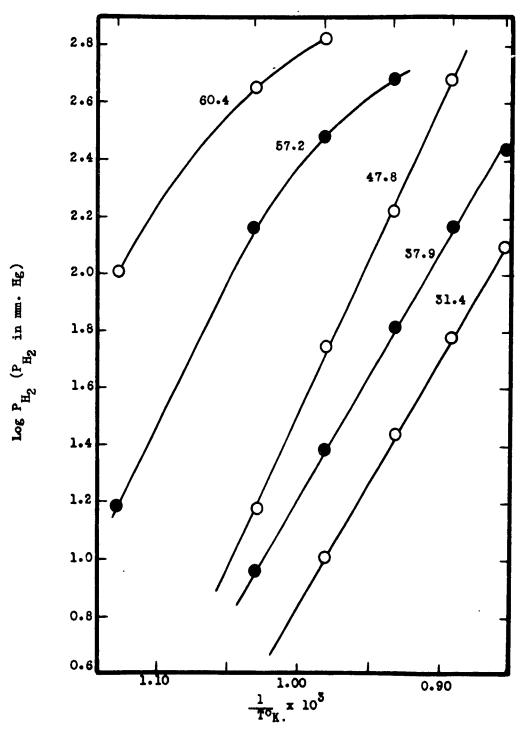


Fig. 17 Pressure-Temperature Curves for an Alloy Containing Initially 6.3 Atom Per Cent Oxygen in the Concentration Range 31.4 to 60.4 Atom Per Cent Hydrogen.

(R = 0.068)

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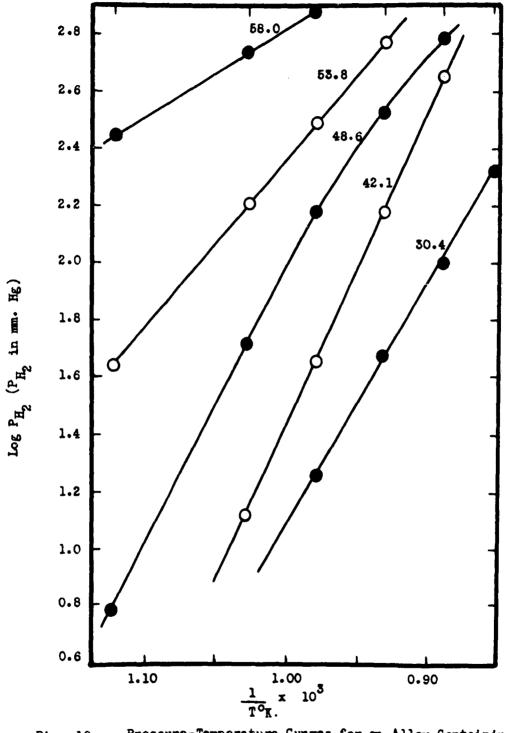


Fig. 18 Pressure-Temperature Curves for an Alloy Containing Initially 10.8 Atom Per Cent Oxygen in the Concentration Range 30.4 to 58.0 Atom Per Cent Hydrogen.

(R = 0.121)

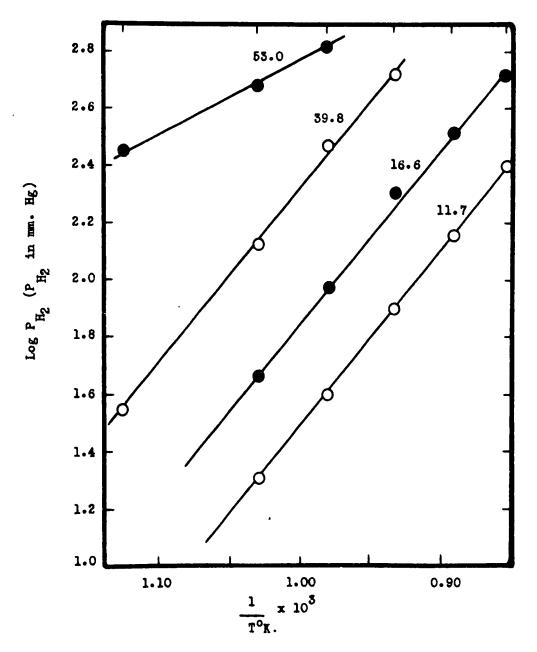


Fig. 19 Pressure-Temperature Curves for an Alloy Containing Initially 18.6 Atom Per Cent Oxygen in the Concentration Range 11.7 to 53.0 Atom Per Cent Hydrogen.
(R = 0.229)

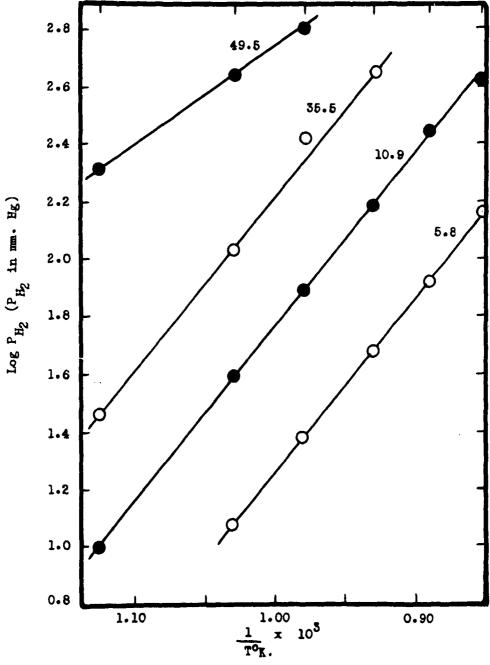


Fig. 20 Pressure-Temperature Curves for an Alloy Containing Initially 24.7 Atom Per Cent Oxygen in the Concentration Range 5.8 to 49.5 Atom Per Cent Hydrogen.

(R = 0.328)

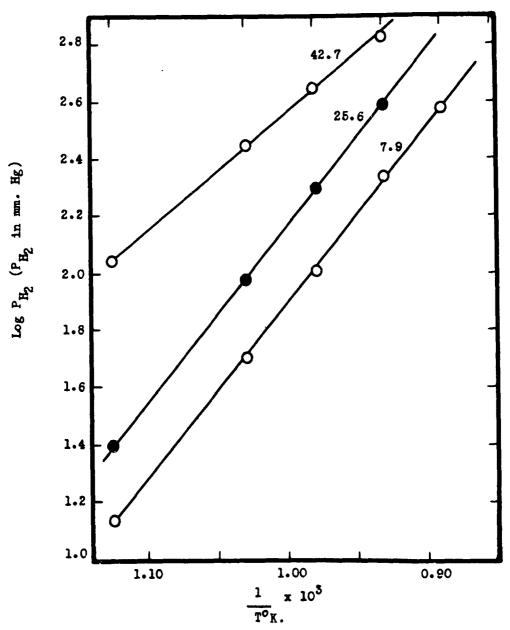


Fig. 21 Pressure-Temperature Curves for an Alloy Containing Initially 29.6 Atom Per Cent Oxygen in the Concentration Range 7.9 to 42.7 Atom Per Cent Hydrogen.

(R = 0.420)

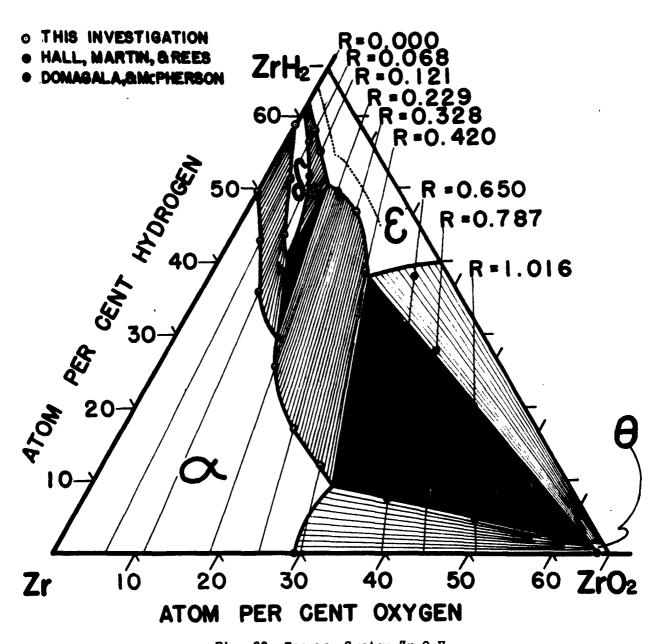


Fig. 22 Ternary System Zr-O-H
Provisional Phase Diagram in Part--Isothermal Cross Section at 750°C.

Two-phase regions are shaded by lines approximating tie-lines. Three-phase regions are drawn in solid black. Extensions of boundaries beyond experimental data is by inference and therefore must be considered provisional. The lines across the diagram leading to R values represent experimental composition traverses as hydrogen was successively added to Zr-O solid solution specimens of constant O/Zr atom ratios, R. Dotted line represents a contour of the variation in composition at constant hydrogen pressure of 760 mm.